

Journal of Structural Geology 27 (2005) 397-408



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# What controls relay ramps and transfer faults within rift zones? Insights from analogue models

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Received 7 May 2004; accepted 8 November 2004 Available online 25 January 2005

#### Abstract

Structures within rift zones exhibit two main types of interaction relevant at the rift scale: relay ramps and transfer faults at high angle to the rift. Analogue experiments have been performed to investigate whether these types of interaction may be affected by differential extension along the rift. In these models, sand (brittle crust analogue) overlies two adjacent silicone (ductile crust analogue) layers with different viscosity, in order to simulate different percentage extension rates ( $\Delta e$ ) along rifts. The experiments show a distinct behaviour as a function of  $\Delta e$ . For  $\Delta e < 21 \pm 3\%$ , extensional structures interact forming relay ramps; for  $\Delta e > 21 \pm 3\%$ , the interaction occurs by means of transfer faults striking subparallel to the extension direction. Experimental data are consistent with the geometries and extension rates of rift zones. Relay ramps characterize narrow rifts and oceanic ridges, where the mean percentage of extension is low (e < 16%). Conversely, transfer faults are usually found in extensional settings (passive margins, wide rifts, back-arc basins) with significant stretching (e>39%), where the rift more likely achieves differential extension  $\Delta e > 21\%$ . © 2005 Elsevier Ltd. All rights reserved.

Keywords: Rift zones; Relay ramps; Transfer faults; Analogue models; Differential extension

# 1. Introduction

Rift zones are commonly segmented at various scales Segments are, however, transient features in the evolution of a rift zone, as they grow and interact and may link to form larger structures (Macdonald and Fox, 1983; Pollard and Aydin, 1988; Dawers and Anders, 1995; Koukouvelas et al., 1999). Therefore, the process of interaction is a necessary step in the evolution of a rift zone over a range of scales.

Two main typologies of interaction between segments (or groups of segments) that appear to be significant at the scale of the considered rift (controlling its shape or continuity) can be identified. The interaction may develop relay ramps or accommodation zones, which consist of broad areas of ductile strain between extensional structures (grabens, normal faults, extensional fractures) usually characterized by arcuate geometries (Fig. 1; e.g. Peacock

et al., 2000a, and references therein). This type of interaction constitutes a common type of 'soft linkage' (Walsh and Watterson, 1991). Conversely, the interaction may develop a transfer fault, which is a subvertical transtensive fault that strikes at high angle and transfers displacement between two adjacent crustal sectors undergoing differential extension (Fig. 1; Gibbs, 1990; Peacock et al., 2000a, and references therein). This type of interaction constitutes a common type of 'hard linkage' (Walsh and Watterson, 1991).

Both types of interaction are found in rift zones. Relay ramps are widespread; they have been described for example in the Rhine Graben (Illies, 1975; Brun et al., 1991), the Rio Grande Rift (Cordell, 1978; Mack and Seager, 1995), the Baikal Rift (Sherman, 1978; Hutchinson et al., 1992), the East African Rift System (EARS) (Morley, 1988; Ebinger, 1989a,b; Ebinger et al., 1989; Morley et al., 1990; Nelson et al., 1992), East Greenland (Larsen, 1988; Peacock et al., 2000b), the British Isles (Peacock and Sanderson, 1991; Huggins et al., 1995; Peacock, 2003), the Aegean Sea (Gawthorpe and Hurst, 1993), the Suez Rift

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Fig. 1. Types of interaction between extensional structures at various rift zones: relay ramps in the East African Rift System (a) (after Ebinger, 1989b) and the Rio Grande graben (b) (after Aldrich, 1986); transfer faults in the Suez Rift (c) (after McClay and Khalil, 1998) and the Atlantic margin of Brazil (d) (after Milani and Davison, 1988); (e) schematic view of relay ramps and transfer faults in extensional domains.

(Moustafa, 1996), the Basin and Range (Anders and Schlische, 1994; Ferrill et al., 1999), the oceanic ridge of Iceland (Acocella et al., 2000) and the Gulf of Thailand (Kornsawan and Morley, 2002). In most of these cases, the relay ramps appear to control the shape or continuity of the rift and therefore they are considered significant at the rift scale (km to  $10^3$  km). Also, the geometric and kinematic features of relay ramps can be consistent over scales from m to  $10^5$  m (Acocella et al., 2000).

Transfer faults were first recognized and described in the Viking Graben, North Sea (Gibbs, 1984). Since then, various authors have reported the occurrence of transfer faults at various extensional settings, such as the Atlantic margin of Brazil (Milani and Davison, 1988), the Atlantic margin of Newark (Schlische, 1992), the Basin and Range (Duenbendorfer and Black, 1992; Martin et al., 1993), the Atlantic margin of Galicia (Boillot et al., 1995), the Atlantic margin of Norway (Dorè et al., 1997; Tsikalas et al., 2001), the Suez Rift (McClay and Khalil, 1998), the Atlantic margin of W Africa (Watts and Stewart, 1998) and the Tyrrhenian (Acocella et al., 1999a) and Japanese (van der

Werff, 2000) back-arc margins. In most of the cases, the transfer faults appear to be relevant at the rift scale (km to  $10^2$  km).

Therefore, while relay ramps are widespread in various extensional settings, transfer faults are usually observed within passive margins, wide rifts and back-arc basins, and are essentially lacking within continental narrow rifts (EARS, Rio Grande Rift, Rhine Graben and Baikal Rift). Also, within readily accessible oceanic extensional domains, such as Iceland, transfer faults are lacking and the dominant type of interaction are relay ramps (Acocella et al., 2000).

To investigate whether the selective occurrence of relay ramps and transfer zones is controlled by the differential extension along the rift, we have performed analogue experiments using sand (brittle crust analogue) and silicone (ductile crust analogue). The experiments simulate crustal blocks undergoing differential extension and show that transfer faults, relevant at the rift scale, can form only with a percentage of differential extension  $\Delta e > 21 \pm 3\%$ . Below this threshold, relay ramps, relevant at the rift scale, occur. This value roughly corresponds in nature to the threshold of

 Table 1

 Model and nature ratios applicable to the present study

Parameter	Model value	Nature value	Model/nature ratio		
Length	0.01 m	10 <sup>4</sup> m	$L^* \sim 10^{-6}$		
Density	$1.2-1.4 \text{ kg/m}^3$	$2.4-2.8 \text{ kg/m}^3$	$\rho^* \sim 0.5$		
Gravity	$9.8 \text{ m/s}^2$	$9.8 \text{ m/s}^2$	$g^* \sim 1$		
Stress $(\sigma = \rho g L)$			$\sigma^* \sim 5 \times 10^{-7}$		
Viscosity ( $\mu = \sigma/\epsilon$ )	$10^4 - 10^5$ Pa s	$10^{20} - 10^{21}$ Pa s	$\mu^* \sim 10^{-16}$		
Strain rate	$5 \times 10^{-6}  \mathrm{s}^{-1}$	$10^{-15}  \mathrm{s}^{-1}$	$\varepsilon^* \sim 5 \times 10^9$		
Time $(t=1/\varepsilon)$			$t^* \sim 2 \times 10^{-10}$		

extension between narrow continental rifts and oceanic ridges, where extension is lower and relay ramps are dominant, and wide rifts, passive margins and back-arc basins, where transfer faults are often more numerous.

#### 2. Experimental procedure

#### 2.1. Scaling and materials

Analogue experiments were constructed to simulate the interaction between adjacent crustal blocks experiencing different extension rates. Scaled models should be geometrically, kinematically and dynamically similar to natural examples (Ramberg, 1981, and references therein). In the experiments, the length ratio L between model and nature is  $L^* = 10^{-6}$  (1 cm in our experiments corresponds to 10 km in nature), the density ratio between rocks and common experimental materials is  $\rho^* \sim 0.5$  and the gravity ratio between model and nature is  $g^*=1$ . The corresponding stress ratio between model and nature is  $\sigma^* = \rho^* g^* z^* \sim 5 \times$  $10^{-7}$  (Table 1). Cohesion c has the dimensions of stress; assuming a Mohr-Coulomb criterion and natural cohesion  $c \sim 10^7$  Pa, a material with  $c \sim 5$  Pa is required to simulate the brittle crust: for this purpose, dry quartz sand, with  $c \sim 0$  Pa, is used; the dry sand has a density  $\sim 1400$  kg/m<sup>3</sup>.



Fig. 2. Sketch of the experimental apparatus. (a) Map view; (b) frontal section view; (c) lateral section view.

Silicone putty with Newtonian behaviour has been used at the base of the sand-pack to simulate the plastically deforming crust. In order to reproduce different amounts of extension across the model, we use two adjacent layers of silicone, with different viscosities of  $7.9 \times 10^4$  Pa s and  $4.5 \times 10^5$  Pa s (Fig. 2b); the silicone has a density ~1310 kg/m<sup>3</sup>.

The following relation applies to Newtonian ductile materials (Benes and Davy, 1996):

$$\sigma_1^* - \sigma_3^* = \mu^* \varepsilon^* \tag{1}$$

where  $\mu^*$  and  $\varepsilon^*$  are the viscosity and the strain rate ratios between model and nature, respectively. Where  $\sigma_1^* - \sigma_3^* \sim 5 \times 10^{-7}$ , the  $\mu^*$  and  $\varepsilon^*$  ratios have to be scaled accordingly (Table 1). Considering the viscosities of silicone and the mean viscosity of the lower crust ( $10^{20}$ –  $10^{21}$  Pa s; Ranalli, 1995),  $\mu^* \sim 10^{-16}$  and, as a result (from Eq. (1)),  $\varepsilon^* \sim 5 \times 10^9$ . For a mean extensional strain rate  $\varepsilon_n \sim 10^{-15}$  s<sup>-1</sup>, commonly found in nature, to have  $\varepsilon^* \sim 5 \times 10^9$ , requires a strain rate  $\varepsilon_m \sim 5 \times 10^{-6}$  s<sup>-1</sup> in the experiments. As  $\varepsilon^* = 1/t^*$ , 1 s in our experiments corresponds to  $2 \times 10^{10}$  s (~634 years) in nature (Table 1).

# 2.2. Set-up

The experimental set-up is characterized by a 1–2-cmthick sand layer, a 1–2-cm-thick silicone layer and a basal plate (Fig. 2). In all the experiments, the silicone layer consists of two adjacent silicone portions with different viscosity (Fig. 2a). The silicone layer is confined on three sides and has a free boundary along the fourth side.

Extension within the model is obtained by means of a moderate tilt  $(4-6^{\circ})$  of the basal plate towards the free boundary (Fig. 2c). The tilt angle has been chosen in order to maintain a consistence between the mean extension rates of the model and nature (Table 1). Such a tilt develops in fact a tangential (parallel to the dip of the tilted plate) component of the gravity force. The gravity force at the base of an experiment characterized by 1.5 cm of sand and 1.5 cm of silicone is:

$$\sigma = (\rho g z)_{\text{sand}} + (\rho g z)_{\text{silicone}} = 398.37 \text{ N}$$

For a basal tilt of  $4^\circ$ , the tangential stress at the base of the experiment is:

Main imposed and observed features in the experiments.  $T_b$ =sand thickness;  $T_d$ =silicone thickness; L=length of model; W=width of model;  $\mu_1$ =viscosity of silicone in one plate;  $\mu_2$ =viscosity of silicone in the other plate;  $\alpha$ =angle of tilting of the basal rigid plate;  $\Delta e$ =percentage of differential extension defining the experimental threshold between relay ramps and transfer faults

Experiment	$T_{\rm b}~({\rm cm})$	$T_{\rm d}~({\rm cm})$	L(cm)	W(cm)	$\mu_1$ (Pa s)	$\mu_2$ (Pa s)	α (°)	$\Delta e$	
TIR 2	1	2	18	28	$4.5 \times 10^{5}$	$7.9 \times 10^{4}$	4	21	
TIR 3	1	2	20	40	$4.5 \times 10^{5}$	$7.9 \times 10^{4}$	6	18	
TIR 4	1.5	1.5	20	40	$4.5 \times 10^{5}$	$2.7 \times 10^{5}$	4	>10	
TIR 5	1.5	1.5	25	40	$4.5 \times 10^{5}$	$7.9 \times 10^{4}$	4	23	
TIR 6	1	2	20	50	$4.5 \times 10^{5}$	$7.9 \times 10^{4}$	4	20	
TIR 7	1.5	2	20	40	$4.5 \times 10^{5}$	$7.9 \times 10^{4}$	4	24	

# $\sigma_{\rm t} = 398.37 \cos 4^\circ = 27.8 \, {\rm N}$

This value represents the tangential component of the gravity force responsible for the flow of silicone along the slope at the beginning of the experiment (Fig. 2c).

The flow of silicone induces the thinning and extension of the overlying sand pack. As the silicone has different viscosities, it undergoes different flow velocities and therefore stretching. Since the amount of extension within the sand is proportional to the amount of stretching of the underlying silicone, this set-up simulates the interaction of adjacent blocks of upper crust with differential extension.

At the end of each experiment, sand was added and levelled onto the surface of the model, to preserve topography. The model was then saturated with water and cut. Six experiments were performed in order to test: the role of the ratio between the thickness of the brittle and ductile materials in the deformation pattern; the silicone viscosity contrast; the slope angle; the lateral dimensions of the models (Table 2).

This apparatus permits us to examine the interaction between structures undergoing differential extension without the control of a rigid basal velocity discontinuity, as commonly used in previous experiments of interacting extensional structures (Courtillot et al., 1974; Elmohandes, 1981; Serra and Nelson, 1988; Naylor et al., 1994; Mauduit and Dauteuil, 1996; Acocella et al., 1999b). The lack of a velocity discontinuity has the important advantage of avoiding the related modifications of the deformation pattern at surface.

# 2.3. Limits and assumptions of the experiments

These experiments study the effect of differential extension within a brittle crust; to achieve this, two silicone putties with different viscosities are used. The use of silicone with different viscosity does not necessarily correspond to the simulation of different types of ductile crust and therefore does not have a specific counterpart in nature. Silicone viscosity is thus varied to simulate a differential extension, not to simulate different types of ductile crust.

The tangential component of the gravity force  $\sigma_t$  responsible for the flow of silicone is not constant during

the experiment. As the sand and silicone layers thin,  $\sigma_t$  decreases:  $\sigma_t$  at the end of the experiment can be estimated at 50% of the initial  $\sigma_t$ . This process is responsible for a decrease in the stretching rate with time, resulting in significant deformation of the model at the earlier stages and moderate deformation at later stages.

Variable extension rates are also found in natural extensional settings and, to a first approximation, assuming a variable extension rate should not limit the applicability of the experiments. Despite the variable extension rate, the total duration (~5 h) and extension (up to ~70%) of the experiments correspond to a realistic duration of rifting (~ $1.1 \times 10^7$  years) and amount of stretching in nature.

The experiments do not take into account any control of pre-existing brittle structures on the development of the interactions, even though these may play an important role. The purpose of the experiments is in fact the study of the type of interaction due to differential extension in the simplest conditions.

Finally, the types of interaction on which this work is focused are those characterizing the continuity and shape of the rift zones; therefore, only those interacting structures that appear relevant at the rift scale, both in nature and the experiments, are considered.

# 3. Experimental results

The evolution of the experiments is here summarized by model TIR 7 ( $T_b = 1.5 \text{ cm}$ ;  $T_d = 2 \text{ cm}$ , where  $T_b$  and  $T_d$  are the sand and the silicone thickness, respectively), with silicone viscosities of  $7.9 \times 10^4$  (left plate in Fig. 3a) and  $4.5 \times 10^5$  Pa s (right plate in Fig. 3a). At t=0' (minutes) the experiment is undeformed (Fig. 3a).

The tilt of the rigid base induces the flowing of silicone and extension in the two plates. At t=60' several depressions form. The plate with more viscous silicone has four regularly spaced depressions bordered by normal faults, whereas the plate with less viscous silicone has six regularly spaced wider depressions bordered by normal faults (Fig. 3b). The lateral termination of these graben-like features along the contact between the adjacent plates is marked by relay ramps or accommodation zones. These consist of a broad deformed area, delimited by normal



Fig. 3. Evolution of experiment TIR 7. Map views of: (a) undeformed stage; (b) experiment at t=60'; (c) experiment at 180'; (d) experiment at 300'. Enlarged section views along: (e) the plate with higher viscosity silicone; (f) the plate with lower viscosity silicone. Arrows in sections indicate the direction of extension; dashed rectangles show examples of horst and graben.



Fig. 4. Time variation of the mean percentage of extension for each plate.

faults, slightly oblique to the extension direction and displaying arcuate shapes.

At t=180' the depressions have become wider and deeper (Fig. 3c). The interaction between the normal faults bordering these grabens now occurs with different modalities. Far from the free boundary, the interaction still occurs through relay ramps (Fig. 3c). Near the free boundary, the normal faults bordering the grabens are interrupted by a set of left-lateral faults subparallel to the extension direction (Fig. 3c). Therefore, the interaction between extensional structures here occurs by means of transfer faults.

At t=300' the experiment does not show significant differences compared with the previous stage, even though the depressions are more accentuated and the strike-slip faults better developed (Fig. 3d).

The section view of the more viscous plate at the end of experiment is shown in Fig. 3e. Its thinning has been achieved through the development of a set of grabens and horst-like structures. The areas of maximum thinning in the brittle part correspond to the areas of rise of the underlying silicone putty.

The section view of the less viscous plate at the end of experiment is shown in Fig. 3f. The horst–graben configuration is here more pronounced, with local elision of the brittle overburden and rise of the silicone. The more severe thinning of this plate is related to the lower viscosity of the silicone.

The mean percentages of extension measured for each plate during the experiment are given by:

$$e_{\rm m} = (L_{\rm m} - L_{\rm i})/L_{\rm i}$$

where  $L_i$  is the initial length of the plate and  $L_m$  is the incremental length of the plate after a given time interval. These percentages represent average values for each plate: at 60' (15 and 32%), 180' (32 and 60%) and 300' (37 and 67%) they show an overall decay in the amount of extension with time (Fig. 4).

The displacement vectors derived from the extension of



Fig. 5. Vectors of displacement of the nodes of the reference grid of experiment TIR 7 at the final stage of deformation (t=300'). The initial position of the nodes coincides with the origin of the arrows.  $e_1$  and  $e_2$  are the percentages of extension related to each couple of nodes (origin of the dotted arrows) at the two sides of the plates contact. The difference ( $e_1 - e_2$ ) between the percentages of extension for each couple of nodes at the sides of the contact gives the percentage of differential extension  $\Delta e$ . The inset highlights the variations of  $\Delta e$  (also represented as shades of grey) along the plate contact.



Fig. 6. Diagram showing the mean spacing of the depressions on the two plates for each experiment. The insets above report the mean spacing relative to all the experiments.

the reference grid at t=300' are shown in Fig. 5. In general, the vectors are longer on the low-viscosity plate and increase towards the free boundary. This implies that the differential extension at the contact between the two plates increases towards the free boundary as well.

Fig. 5 also shows the percentages of extension ( $e_1$  and  $e_2$ ) related to each couple of nodes (origin of the dotted arrows) at the border between the plate contact; these percentages are given, for each node, by:

 $e = (L_{\rm f} - L_{\rm i})/L_{\rm i}$ 

where  $L_i$  is the initial length of the model and  $L_f$  is the incremental length of the node plus the initial length of the model  $L_i$ . The difference  $(e_1 - e_2)$  between the percentages of extension for each couple of nodes at the sides of the

contact gives the local percentage of differential extension  $\Delta e$  (inset in Fig. 5).

The shades of grey in Fig. 5 show that the presence of transfer faults (grey lines) connecting the two extending plates is limited to differential extension  $\Delta e > 24\%$ ; below this threshold, the extending plates are connected by relay ramps. The differential extension of 24% represents therefore, in this experiment, the threshold between two types of interaction between extensional structures, characterized by relay ramps and transfer faults.

The remaining experiments, with the exception of TIR 4, showed an overall deformation pattern similar to TIR 7, despite their different model attributes (Table 2). The similar deformation pattern was produced during the development of relay ramps far from the free boundary and transfer faults near to the boundary. The values of differential extension associated with the presence of transfer faults are similar to TIR 7, giving a mean threshold of  $21\pm3\%$  (Table 2). These data show that the presence of transfer faults in all the experiments is restricted, at any time during the evolution of an experiment, to a mean differential extension >21%; at lower values, relay ramps occur.

The depressions in both plates in all the experiments are regularly spaced. Their mean spacing, related to each plate and experiment, is shown in Fig. 6. The depressions on the lower viscosity plate have a slightly larger mean spacing (S=6.7) with regard to those on the higher viscosity plate (S=4.8).

The experimental transfer faults are usually arranged in subparallel segments. Their evolution is shown through different stages of experiment TIR 3 (Fig. 7); transfer faults usually grow in length through the linkage of en-échelon segments. Their propagation is mainly away from the free



Fig. 7. Evolution of a transfer fault in experiment TIR 3.

Table 3

Estimates of $\beta$ (and related references) for rift zones on Earth,	Venus and Mars.	The references	describing the ty	pe of interaction (	r.r.=relay	ramps, t.f. $=$
transfer faults) are given in the introduction						

Rift type	Reference	Stretching $\beta$	Interaction type	
Narrow rifts				
Rio Grande	(Golombek et al., 1983)	1.08	r.r.	
Rhine graben	(Illies, 1979; Villemin et al., 1986)	$1.13 \pm 0.03$	r.r.	
East African Rift System	(Ebinger, 1989a; Prodehl et al., 1997)	1.15	r.r.	
Baikal	(Agar and Klitgord, 1995)	1.16	r.r.	
Oceanic ridges				
Iceland	(Forslund and Gudmundsson, 1991; Dauteuil et al., 2001)	<1.06	r.r.	
Extra-terrestrial narrow rifts				
Valles Marineris (Mars)	(Mege and Masson, 1996)	1.06	r.r.	
Beta Regio (Venus)	(Foster and Nimmo, 1996)	$1.1 \pm 0.1$	r.r.	
Wide rifts				
Mojave Desert	(Martin et al., 1993)	1.5	t.f.+r.r.	
Basin and Range	(Wernicke, 1985)	>1.5	t.f.+r.r.	
North Sea	(Latin and White, 1990)	2.5	t.f.+r.r.	
Passive margins				
Brazil	(Milani and Davison, 1988)	1.39	t.f.+r.r.	
Newark	(Schlische, 1992)	1.45	t.f.+r.r.	
Suez	(Angelier, 1985)	1.45	t.f.+r.r.	
NE Atlantic	(Tsikalas et al., 2001)	2.2	t.f.+r.r.	
Galicia	(Boillot et al., 1995)	>2.75	t.f.+r.r.	
W Africa	(Watts and Stewart, 1998)	>3	t.f.+r.r.	
Back-arc basins				
Tyrrhenian margin	(Faccenna et al., 1997)	>1.45	t.f.+r.r.	
Japanese margin	(Jolivet et al., 1994)	1.8	t.f.+r.r.	

boundary, because the percentage of differential extension is greatest at the free boundary and increases with time at all points along the transfer zone.

TIR 4 was the only experiment where transfer faults did not develop and was characterized by a lower viscosity contrast between the two silicone layers (viscosities of  $1.7 \times 10^5$  and  $4.5 \times 10^5$  Pa s). As a consequence, their maximum differential extension was ~10% and, consistently with the above results, the interaction between extensional structures in the two plates was characterized only by relay ramps.

## 4. Discussion and conclusions

#### 4.1. Interpretation of the experiments

The tilt of the basal plate induces the downward flow of the silicone, causing extension of the model. At the very beginning, the low amount of extension within a restricted area is responsible for a deformation pattern interpretable as similar to the one of narrow rifts. At later stages, the higher amount of extension within a wider area is responsible for a deformation pattern more similar to the one of wide rifts. The rate of extension shows an overall decrease with time (Fig. 4), consistently with the decrease of  $\sigma_t$ .

The different viscosities of silicone at the base of the models are responsible for the differential extension. The experiments are all consistent with a differential extension associated with graben-like structures; these are, on the less viscous plate, more numerous, wider and deeper. The grabens show moderate variations in their spacing (Fig. 6), controlled by the interplay between the variations in  $T_{\rm b}$ ,  $T_{\rm d}$ and the viscosity of silicone. In particular, the spacing Sbetween instabilities responsible for the initiation of stretching of competent materials with thickness  $T_{\rm b}$  follows the relationship  $S \sim 4T_{\rm b}$  (Ricard and Froidevaux, 1986). Also, the instabilities developed during extension lead to the thinning of the sand overburden and the consequent rise of silicone; the lower the viscosity and the higher the thickness of silicone, the easier it will rise (Brun, 1999). A higher amount of risen silicone results in the higher spacing of the depressions (Fig. 3f). Similar processes have been used to study different modalities of continental extension, assuming therefore a general significance (Brun, 1999, and references therein).

Since the silicone nearer to the free boundary moves further (Fig. 5), the differential extension  $\Delta e$  ( $\Delta e = e_1 - e_2$ )

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increases towards the free boundary. As a result, two types of interaction develop due to the gradient of differential extension.

The type of interaction is a function of the percentage of differential extension between the two plates  $\Delta e$ . When  $\Delta e < 21 \pm 3\%$ , the interaction occurs through relay ramps or accommodation zones. Their evolution can be characterized by the following stages (Fig. 3): (a) lateral propagation of the faults; (b) linkage with the adjacent structure. Their overall evolution is therefore similar to that of the natural prototypes (e.g. Acocella et al., 2000, and references therein). When  $\Delta e > 21 \pm 3\%$ , the interaction occurs through transfer faults. The experiments display the growth of the transfer faults, due to the linkage of smaller enéchelon strike-slip segments (Fig. 7): because of their geometries and kinematics, these segments can be interpreted as Reidel systems *R* of the growing transfer fault.

The experiments suggest that relay ramps, usually characterized by a minor component of strike-slip, accommodate minor differential displacements between two adjacent extensional structures. Conversely, larger differential displacements can be only accommodated by predominant strike-slip systems parallel to the extension direction, such as transfer faults.

The very moderate scatter  $(\pm 3\%)$  of the experimental threshold suggests that geometric parameters (such as the ratio between the thickness of the brittle and ductile materials, the slope angle and the lateral dimensions of the models) do not significantly affect the overall modalities of deformation, the types of interaction and their threshold value. The 3% scatter is moderate also when compared with the variations in the spacing (approximately  $\pm 20\%$ ; Fig. 6) of the depressions; this suggests that the development of the type of interaction is independent from the overall configuration (frequency or spacing) of the extensional structures. Moreover, the fact that the two types of interaction occur above the contact between the silicone layers confirms that the latter does not influence the type of deformation, depending entirely on the amount of differential extension.

# 4.2. Comparison with nature

A qualitative comparison between our experiments and nature shows that the analogue faults share close geometric and kinematic similarities with those observed in extensional settings Also, the overall deformation pattern is consistent with the ones observed at various rifts, such as the Basin and Range, the Aegean Sea or the Tyrrhenian area (Duenbendorfer and Black, 1992; Gawthorpe and Hurst, 1993; Martin et al., 1993; Acocella and Funiciello, 2004).

A quantitative comparison is limited by the insufficient knowledge of the percentages of differential extension between crustal portions in rift zones. This constitutes a significant limitation for the complete and rigorous application of the results to nature. In fact, what is usually known about a rift zone is its overall stretching factor  $\beta$ , where  $e = (\beta - 1) \times 100$ , rather than the differential extension between its portions.

The  $\beta$  estimates for various extensional settings and the related references are shown in Table 3; the references describing the type of interaction observed in each setting are listed in the introduction. Table 3 shows that the narrow rifts (EARS, Baikal Rift, Rhine Graben, Rio Grande Rift) are characterized by a stretching factor  $\beta < 1.16$ . In contrast, the wide rifts (Basin and Range, North Sea), back-arc basins (Tyrrhenian Sea, Japan Sea) and passive margins at various stages of evolution (Suez Rift, Atlantic margins of Brazil, W Africa, Newark, Galicia and Norway) are characterized by higher stretching factors, where  $\beta > 1.39$ .

These data suggest that at narrow rifts, characterized by a limited amount of stretching ( $\beta < 1.16$ ), significant (relevant at the rift scale) transfer faults cannot usually form. In fact, as the experiments suggest that transfer faults occur for  $\Delta e > 21\%$ , a differential extension within the narrow rift requires, to form transfer faults, stretching values higher  $(\beta > 1.21)$  than those measured. As a result, at narrow rifts relay ramps are the commonly observed type of interaction between extensional structures (Illies, 1975; Cordell, 1978; Sherman, 1978; Morley, 1988; Ebinger, 1989a,b; Brun et al., 1991; Hutchinson et al., 1992; Mack and Seager, 1995) and significant transfer faults are lacking. The formation of transfer faults for  $\Delta e < 21\%$  may indeed be possible, provided there is a presence of pre-existing structures, subparallel to the extension direction, along the boundary between extending crustal portions. In this case, their reactivation may form transfer faults. Evidence for reactivated transfer faults is found in the Apennines of central Italy, where pre-existing NE-SW structures are reactivated as transtensive under NW-SE regional extension, with  $\Delta e < 21\%$  (Acocella and Funiciello, 2004). Also, transfer-like features in the EARS have been interpreted as due to the reactivation of pre-existing structures (Rosendahl, 1987).

Conversely, on wide rifts, back-arc basins and passive margins, where  $\beta > 1.39$ , the kinematic conditions required to develop transfer faults can be fully met. Where the percentage of differential extension remains at  $\Delta e < 21\%$ , relay ramps continue to form. Nevertheless, because of the larger extension involved, a non-uniform extension may locally allow  $\Delta e > 21\%$ , developing transfer faults; this could be more easily achieved in those settings characterized by the highest  $\beta$  (Table 3). As a result, both transfer faults (Gibbs, 1984; Milani and Davison, 1988; Schlische, 1992; Martin et al., 1993; Boillot et al., 1995; Clemson et al., 1997; Dorè et al., 1997; McClay and Khalil, 1998; Watts and Stewart, 1998; Acocella et al., 1999a; van der Werff, 2000; Tsikalas et al., 2001) and relay ramps (Larsen, 1988; Gawthorpe and Hurst, 1993; Anders and Schlische, 1994; Moustafa, 1996; Ferrill et al., 1999; Peacock et al., 2000b) are commonly observed as types of interaction between extensional structures where  $\beta > 1.21$  (Table 3).



Fig. 8. Types of interaction within rift zones on Earth, Venus and Mars as a function of the stretching factor  $\beta$ . Relay ramps are widespread and independent from  $\beta$ . Transfer faults are observed only where  $\beta > 1.39$ . The experimental threshold (e=21% or  $\beta=1.21$ ) coincides with the separation between domains where relay ramps (continental narrow rifts, oceanic ridges and extra-terrestrial narrow rifts) and relay ramps and transfer faults (wide rifts, passive margins, back-arc basins) have been observed.

These considerations show an overall consistency between the experimental data and the continental extensional domains. As far as oceanic extensional domains are concerned, the best-studied ridge is possibly the Icelandic Ridge. Calculations of its crustal dilation suggest a maximum  $\beta$ =1.06 (Forslund and Gudmundsson, 1991; Dauteuil et al., 2001) and therefore transfer faults should be inhibited. This is in agreement with the fact that transfer faults are not observed within the ridge of Iceland and that the dominant type of interaction are relay ramps (Acocella et al., 2000).

Rift zones are also found on other terrestrial planets, characterized by a rigid lithosphere. Among these, the bestknown cases are Beta Regio (Venus) and Valles Marineris (Mars), both consisting of narrow rifts, with overall deformation patterns similar to the rifts on Earth (Frey, 1979; Foster and Nimmo, 1996; Anderson et al., 2001). These planets exhibit gravity forces and crustal thickness that are different to those on Earth. Even though the performed experiments have not been specifically built to take into account these variations, it is nevertheless interesting to consider that these rift zones are likely associated with a maximum  $\beta = 1.1$  (Foster and Nimmo, 1996; Mege and Masson, 1996) and lack of transfer faultlike structures. Therefore, their overall geometry and kinematics may be consistent with the experimental results and the Earth analogues.

The comparison between the experimental data and the considered continental and oceanic rifts on Earth and selected examples form terrestrial planets suggests a consistency in the type of interaction between extensional structures within rift zones (Fig. 8). This consistency consists of the widespread presence of relay ramps within rift zones characterized by moderate stretching factors ( $\beta <$ 

1.16) and both of transfer faults and relay ramps within rift zones characterized by higher stretching factors ( $\beta > 1.39$ ). In the latter case, it is proposed that the amount of differential stretching  $\Delta e$  between adjacent extensional structures will locally determine the occurrence of transfer faults ( $\Delta e > 21\%$ ) or relay ramps ( $\Delta e < 21\%$ ).

#### Acknowledgements

C. Faccenna is acknowledged for helping in the set-up of the experiments and a critical read of the manuscript. N. D'Agostino kindly provided Fig. 5. Suggestions from J.P. Brun, A. Nicol and A.E. Clifton helped to significantly improve the work. Work partly financed with GNV (Campi Flegrei Project) Funds.

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